

Precision Electro-Weak Parameters from AdS-CFT

[hep-ph/0608241](#) (40 pp.)

[hep-ph/0609104](#) (4 pp.)

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...in three Acts.

- Introductory hand-waving.
- A well defined model and a detailed calculation.
- Interpreting the result.



Conclusions

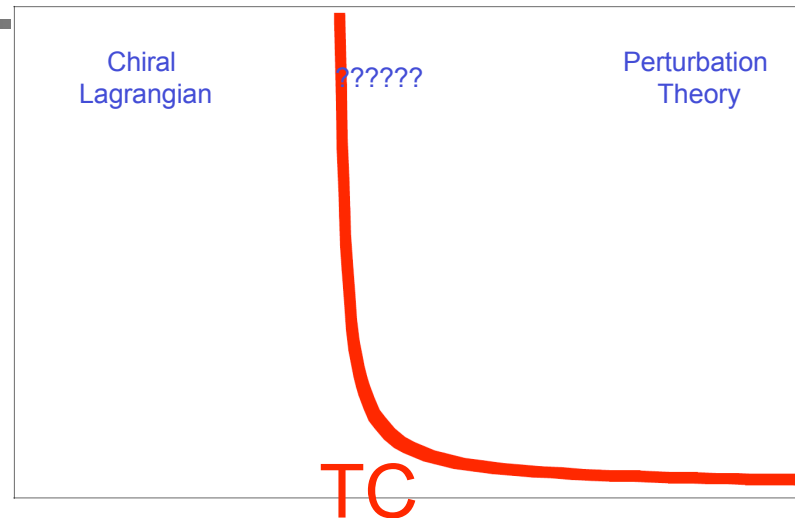
- AdS-CFT correspondence useful tool.
 - Non-perturbative effects are (=can be) huge.
 - Walking-TC compatible with data.
 - Experimental bounds are N_c -independent.
 - Spin-1 resonances at 2 TeV.
 - Degenerate spectrum of spin-1 states.
 - No light scalar (=very broad Higgs at 1-2 TeV?)
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- Model building to be done. Can be done.
 - LHC phenomenology to be studied. Can be studied.



A Dead Horse

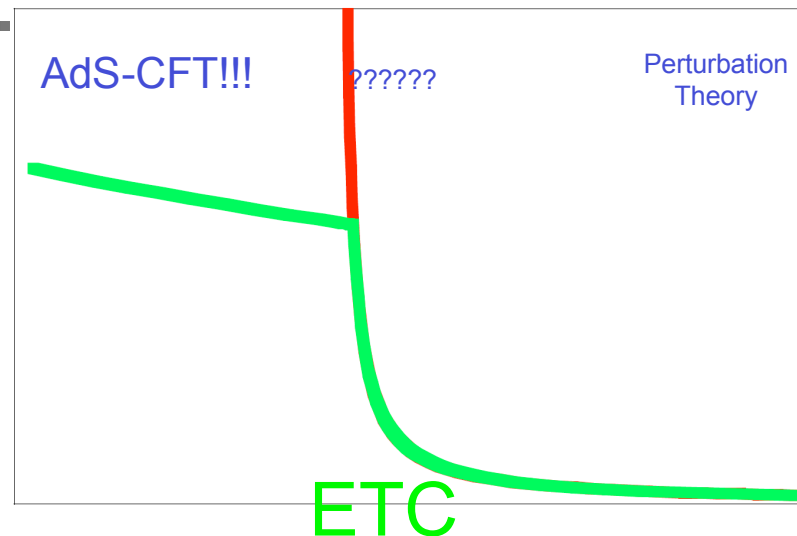
- After LEP, SLAC and Tevatron, Technicolor (naïf version of) dismissed, because it does too much:
 - S too big.
 - T too big.
 - Top mass too small.
 - Too many PNGB's.
 - Too much FCNC.
 - Incomprehensible CKM.
- Too difficult to compute something.
- Too difficult to build a model.

The “Why”



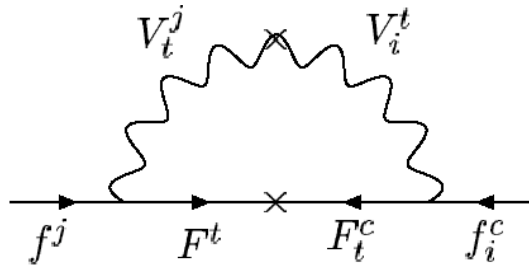
- ONE dynamical scale $TC \sim EWSB$.
- Higher-Order operators unsuppressed at electro-weak scale (Little Hierarchy, S, T, FCNC...)
 - Computational nightmare at electroweak scale
- Only good: NO big hierarchy problem (conformal symmetry at weak coupling)

The Solution



- TWO (maybe more...) dynamical scales $ETC \gg TC \sim EWSB$.
- Higher-Order operators suppressed by large scale (S, T, FCNC...)
 - Conformal Symmetry below ETC: little hierarchy solved!
- Computational nightmare at ETC scale $\sim 5-10$ TeV: BUT who cares!!!!
 - Conformal Symmetry at Large Coupling: Large anomalous dimensions, a new computational tool is need. AdS-CFT!!!!

The Top Mass



$$M_{\text{top}} \sim \left(\frac{g}{\sqrt{2}} \right)^2 \eta \frac{\langle Q^{tT} C U_t^c \rangle}{M_{ETC}^2}$$

$$= \frac{8\pi}{3a^2} \frac{\Lambda_F^3}{\Lambda_{ETC}^2} \eta \lesssim \frac{8\pi}{3} \frac{\Lambda_F^2}{\Lambda_{ETC}}$$

- If the chiral condensate has dimension $d=3$, the top mass is ways too small.
- In a CFT at large coupling, there is no reason to think the anomalous dimensions be perturbative. $d < 3$ reasonable.
- For $d < 3$ top mass parametrically enhanced. If $d=2$ and $ETC \sim 4-5$ TeV, estimates not parametrically small (maybe topcolor ?) .



Precision Parameters

- Defined in terms of the polarizations:
 $\hat{S} \equiv \frac{g_4}{g'_4} \pi'_{WB}(0) ,$
 $\mathcal{L} = \frac{P_{\mu\nu}}{2} A_i^\mu \pi_{ij}(q^2) A_j^\nu + g_4^a J_{a\mu} A_a^\mu$
 $\hat{T} \equiv \frac{1}{M_W^2} (\pi_{WW}(0) - \pi_+(0)) ,$
- Tight Experimental Constraints (mH~800 GeV?):

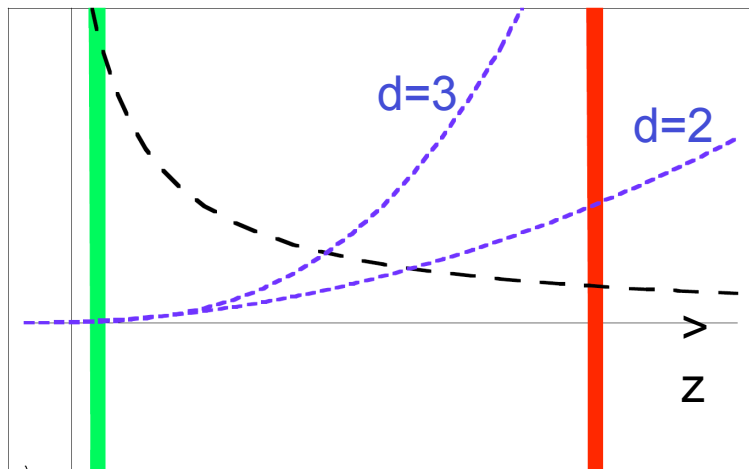
$$\hat{S}_{exp} = (-0.9 \pm 3.9) \times 10^{-3} ,$$

$$\hat{T}_{exp} = (2.0 \pm 3.0) \times 10^{-3} ,$$

- Custodial Symmetry.
- No Non-Perturbative Estimate for S (as of July 2006).
- Perturbative Estimates are BIG (unless $N_c N_d < 8$)

$$\hat{S}_p = \frac{\alpha}{4 \sin^2 \theta_W} \frac{N_c N_d}{6\pi}$$

- ...but why should we trust this?
- ...what is the error?



- $$ds^2 = \left(\frac{L}{z}\right)^2 \left(\eta_{\mu\nu} dx^\mu dx^\nu - dz^2\right)$$
- **Boundaries:** $L_0 < z < L_1$
 - **Consistency:** $L_0 > L$

$$ds^2 = \left(\frac{L}{z}\right)^2 \left(\eta_{\mu\nu} dx^\mu dx^\nu - dz^2\right)$$

$$L_0 < z < L_1$$

$$L_0 > L$$



The Model: Action

$$\mathcal{S}_5 = \int d^4x \int_{L_0}^{L_1} dz \sqrt{G} \left[(G^{MN} (D_M \Phi)^\dagger D_N \Phi - M^2 |\Phi|^2) \right. \\ \left. \left(-\frac{1}{2} \text{Tr} (W_{MN} W_{RS}) - \frac{1}{4} B_{MN} B_{RS} \right) G^{MR} G^{NS} \right] \quad \begin{array}{l} \Phi \sim (2, 1/2) \\ SU(2)_L \times U(1)_Y \end{array}$$

$$\mathcal{S}_4 = \int d^4x \int_{L_0}^{L_1} dz \sqrt{G} \left[\delta(z - L_0) G^{\mu\rho} G^{\nu\sigma} \right. \\ \left[-\frac{1}{2} D \text{Tr} [W_{\mu\nu} W_{\rho\sigma}] - \frac{1}{4} D B_{\mu\nu} B_{\rho\sigma} \right] \\ \left. - \delta(z - L_i) 2\lambda_i \left(|\Phi|^2 - \frac{v_i^2}{2} \right)^2 \right]$$

- Kinetic boundary terms needed for renormalization.
- Boundary terms introduce spontaneous EWSB.



EWSB

- Bulk VEV for Higgs: $\langle \Phi \rangle = \frac{v(z)}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$
- Bulk Equations: $\partial_z \left(\frac{L^3}{z^3} \partial_z v \right) - \frac{L^5}{z^5} M^2 v = 0$
- Solution: $M^2 = -4/L^2$
 $v(z) = Az^2 + Bz^2 \log(z/L)$
- Boundary terms: $\lambda_i \rightarrow +\infty$
 $v(L_0) = v_0,$
 $v(L_1) = v_1,$
- Finally d=2: $\frac{v_0}{L_0^2} = \frac{v_1}{L_1^2}$
 $v(z) = \frac{v_1}{L_1^2} z^2 = \frac{v_0}{L_0^2} z^2$



Electro-Weak Phenomenology

- Define:

$$V^M \equiv \frac{g'W_3^M + gB^M}{\sqrt{g^2 + g'^2}}$$

$$A^M \equiv \frac{gW_3^M - g'B^M}{\sqrt{g^2 + g'^2}}$$

- Bulk Equations:

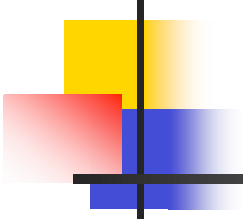
$$A^\mu(q, z) \equiv A^\mu(q)v_Z(z, q)$$

$$\partial_z \frac{L}{z} \partial_z v_i - \mu_i^4 L z v_i = -q^2 \frac{L}{z} v_i$$

- Where:

$$\mu_W^4 = 1/4 g^2 v_0^2 / L^2$$

$$\mu_Z^4 = 1/4 (g^2 + g'^2) v_0^2 / L^2.$$



$$\mathcal{L} = \frac{P_{\mu\nu}}{2} A_i^\mu \pi_{ij}(q^2) A_j^\nu + g_4^a J_{a\mu} A_a^\mu$$

■ Polarizations from UV-boundary Action

$$\frac{\pi_+}{\mathcal{N}^2} = Dq^2 + \frac{\partial_z v_W}{v_W}(q^2, L_0),$$

$$\frac{\pi_{BB}}{\mathcal{N}^2} = Dq^2 + \frac{g^2}{g^2 + g'^2} \frac{\partial_z v_v}{v_v}(q^2, L_0) + \frac{g'^2}{g^2 + g'^2} \frac{\partial_z v_Z}{v_Z}(q^2, L_0),$$

$$\frac{\pi_{WB}}{\mathcal{N}^2} = \frac{gg'}{g^2 + g'^2} \left(\frac{\partial_z v_v}{v_v}(q^2, L_0) - \frac{\partial_z v_Z}{v_Z}(q^2, L_0) \right),$$

$$\frac{\pi_{WW}}{\mathcal{N}^2} = Dq^2 + \frac{g'^2}{g^2 + g'^2} \frac{\partial_z v_v}{v_v}(q^2, L_0) + \frac{g^2}{g^2 + g'^2} \frac{\partial_z v_Z}{v_Z}(q^2, L_0),$$



Regularization

- Taking: $L_0 \rightarrow L$
- Expanding for: $L_0 \rightarrow 0$

$$\frac{\partial_z v_v}{v_v}(q^2, L_0) = q^2 L_0 \left(\frac{\pi}{2} \frac{Y_0(qL_1)}{J_0(qL_1)} - \left(\gamma_E + \ln \frac{qL_0}{2} \right) \right)$$

$$\frac{\partial_z v_Z}{v_Z}(q^2, L_0) = L_0 \left\{ \mu_Z^2 - q^2 \left[\gamma_E + \ln(\mu_Z L_0) + \frac{1}{2} \psi \left(-\frac{q^2}{4\mu_Z^2} \right) - \frac{c_2}{2c_1} \Gamma \left(-\frac{q^2}{4\mu_Z^2} \right) \right] \right\}$$

- From Neumann at IR:

$$c_1 = 2L \left(-1 + \frac{q^2}{4\mu_Z^2}, \mu_Z^2 L_1^2 \right) + L \left(\frac{q^2}{4\mu_Z^2}, -1, \mu_Z^2 L_1^2 \right),$$

$$c_2 = -U \left(-\frac{q^2}{4\mu_Z^2}, 0, \mu_Z^2 L_1^2 \right) + \frac{q^2}{2\mu_Z^2} U \left(1 - \frac{q^2}{4\mu_Z^2}, 1, \mu_Z^2 L_1^2 \right)$$



Renormalization

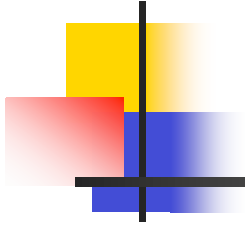
- Define, at finite UV cut-off:

$$D = L_0 \left(\ln \frac{L_0}{L_1} + \frac{1}{\varepsilon^2} \right)$$

$$\mathcal{N}^2 = \varepsilon^2 / L_0$$

- Cut-off dependence disappears, take the limit UV cut-off \rightarrow Infinity. Gauge coupling kept fixed:

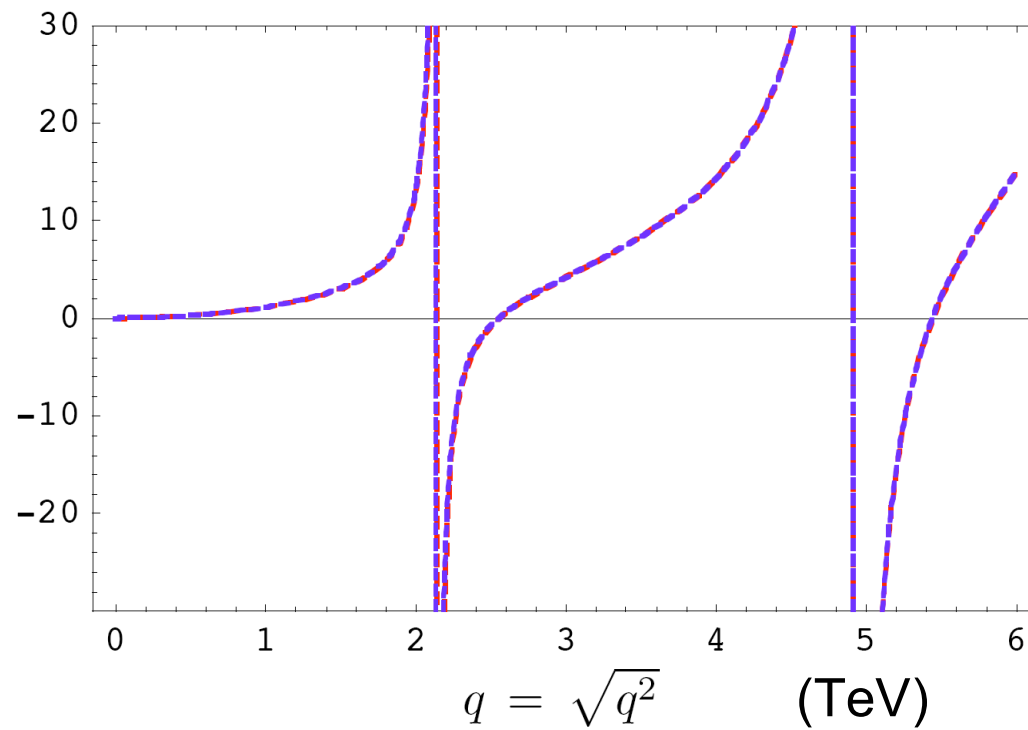
$$g_4^{(\prime)2} = \varepsilon^2 g^{(\prime)2} / L$$

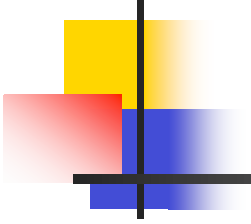


Polarizations

$$g/\sqrt{L} \sim 1.3$$

$$M_{\rho^0} \simeq 2.5 \text{ TeV}$$

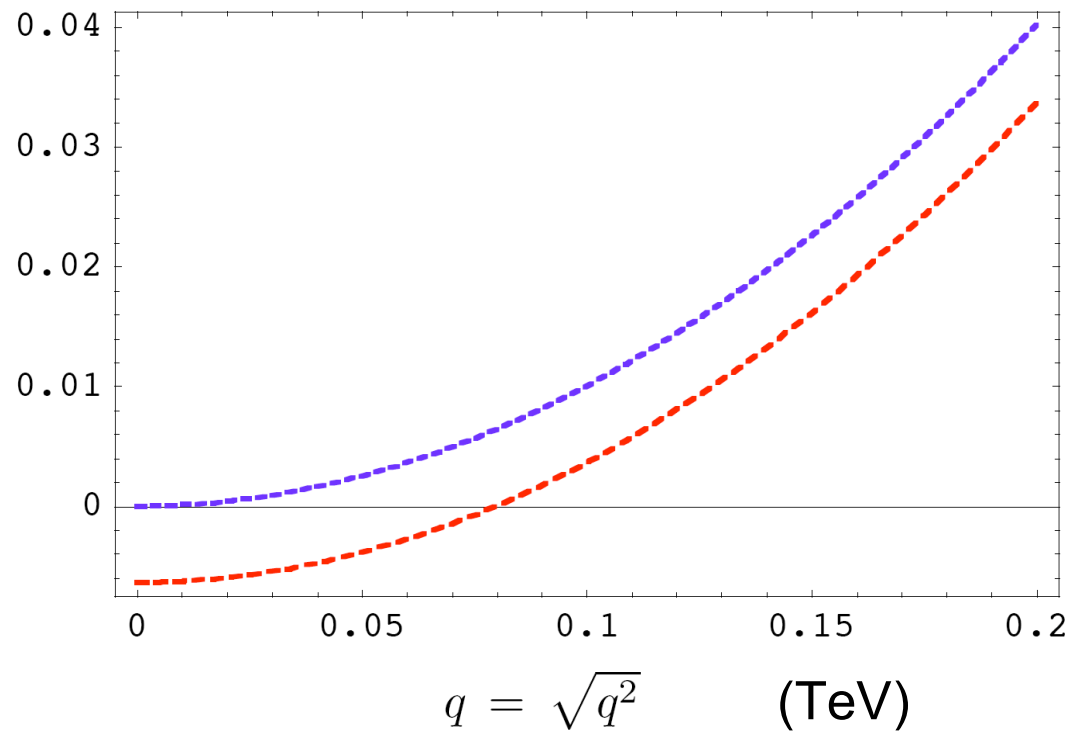




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Phenomenology

- Assume: $\mu_Z^2 L_1^2 \ll 1$
- Spectrum: $M_{\rho^0} = k/L_1 \quad k \in [2.4, 4.7]$

$$M_W^2 \simeq \varepsilon^2 \left(\mu_W^2 \tanh \frac{\mu_W^2 L_1^2}{2} \right) \simeq \frac{1}{2} \varepsilon^2 \mu_W^4 L_1^2,$$

$$M_Z^2 \simeq (g^2 + g'^2)/g^2 M_W^2$$

- EW precision observables:

$$\begin{aligned} \hat{T} &= \frac{\varepsilon^2}{M_W^2} \left(\mu_W^2 \tanh \frac{\mu_W^2 L_1^2}{2} - \frac{\mu_W^4}{\mu_Z^2} \tanh \frac{\mu_Z^2 L_1^2}{2} \right) \\ &\simeq \frac{\varepsilon^2}{M_W^2} \frac{\mu_W^4 L_1^6}{24} (\mu_Z^4 - \mu_W^4) \end{aligned}$$

$$\hat{S} = \varepsilon^2 \frac{1}{2e} \mu_W^4 L_1^4$$



Experimental Bounds

- Experiment: $\hat{S}_{exp} = (-0.9 \pm 3.9) \times 10^{-3},$

$$\hat{T}_{exp} = (2.0 \pm 3.0) \times 10^{-3},$$

- Theory: $\hat{S} \simeq \frac{1}{e} M_W^2 L_1^2 = \frac{k^2}{e} \frac{M_W^2}{M_{\rho^0}^2},$

$$\hat{T} = \frac{M_Z^2 - M_W^2}{6\varepsilon^2} L_1^2 = \frac{k^2}{6\varepsilon^2} \frac{M_Z^2 - M_W^2}{M_{\rho^0}^2}$$

- Bounds:

$$\frac{1}{L_1} > \frac{M_W}{\sqrt{e\hat{S}_{\max}}} = 890 \text{ GeV}$$

- Techni-rho mass:

$$\varepsilon > 1/2 \quad (g/\sqrt{L} < 1.3)$$

$$k(\varepsilon = 1/2) \simeq 2.8.$$

$$M_{\rho^0} \simeq 2.5 \text{ TeV}$$



Fine-Tuning?

- Bounds evaded by: $\mu_Z^2 L_1^2 \ll 1$
- Look back at regularized theory:

$$M_W^2 = \frac{1}{8} \varepsilon^2 g^2 v_1^2 \left(\frac{L_0}{L_1} \right)^2 = \frac{1}{4} g_4^2 \eta^2$$
$$\eta^2 = L \frac{v_1^2}{2} \left(\frac{L_0}{L_1} \right)^2 = \frac{1}{\sqrt{2} G_F} \simeq (246 \text{ GeV})^2$$

- Translation: $\frac{v_1^2 L L_0^2}{2} = \eta^2 L_1^2 < \left(\frac{1}{3.6} \right)^2$
- Is it NATURAL?



Some Estimates

- “Natural” value: $v_1 \simeq \frac{2.4}{gL_1}$
- From QCD... $\sqrt{2}g_\rho f_\pi = M_\rho$
 $g_\rho = g/\sqrt{L}$
- ..and large q ... $L/g^2 = N_c/12\pi^2$
 $g_\rho \simeq 6$
- Conclusion: $\frac{L}{g^2} \frac{L_0^2}{L_1^2} < \left(\frac{1}{6}\right)^2$

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- “Natural” value: $v_1 \simeq \frac{2.4}{gL_1}$
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- ..and large q... $L/g^2 = N_c/12\pi^2$
 $g_\rho \simeq 6$ **NEW NON-PERTURBATIVE**
- Conclusion:
PERTURBATIVE RESULT $\frac{L}{g^2} \frac{L_0^2}{L_1^2} < \left(\frac{1}{6}\right)^2$



More Estimates

- UV cut-off:

$$g/\sqrt{L} \sim 1.3$$

$$1/L_0 \sim 6/L_1 \sim 5.3 \text{ TeV}$$

- Localized Top:

$$-\delta(z-L_0)\tilde{y}_u\bar{q}_L\tilde{\Phi}u_R$$

$$\frac{y_u}{\sqrt{L}} = \frac{L_1}{\sqrt{2}L_0}$$

$$y_u/\sqrt{L} \sim 4$$

- Perturbative:

$$N_d \sim 2N_T \quad N_T \sim 8.$$

$$\hat{S}_p = \frac{\alpha}{4\sin^2\theta_W} \frac{N_d N_T}{6\pi} \sim 0.06 \quad \text{vs.} \quad \hat{S} \simeq 0.003$$



Systematic Errors

- Large N: 5% ?
- Model Dependences: 50% ??
- Departure from ADS5: 50% ??
- Higher order operators: 50% ??

VS.

- Perturbative Estimate: 2000% !!!!!!!



What's next?

- LHC-phenomenology: production cross-sections and decay rates.
- LHC-phenomenology: where is the Higgs?
- Fine-tuning study: stabilization a` la GW?.
- Fermion model-building: hierarchies in mass? CKM? FCNC?
- Generalizations: are T and S always positive? Is there a simple formula for general d ? What about departures from AdS5?



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